Título The time derivative in the mass transport equation in a Drift-flux model

Tipo de Producto Ponencia resumen

Autores Axel Larreteguy, Pablo Caron y Francisco Barceló

Publicado en: XXIII Congreso de Métodos Numéricos y sus Aplicaciones, ENIEF 2017. Universidad Nacional de La Plata.

Código del Proyecto y Título del Proyecto

P17T03 - Simulación de Flujos Multifásicos en la Industria del Petróleo

Responsable del Proyecto

Axel Larreteguy

Línea

Fluidodinámica Computacional (FC)

Área Temática

Modelado y Simulación Computacional (MYS)

Fecha

Noviembre 2017





The time derivative in the mass transport equation in a drift-flux model

Axel Larreteguy, Pablo Caron, Francisco Barceló

Instituto de Tecnología - Universidad Argentina de la Empresa

ENIEF 2017, 5-10 de noviembre, UNLP, La Plata, Argentina





Larreteguy - Caron - Barceló

Time derivative in drift-flux method

Table of Contents

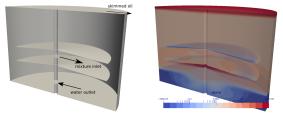
UADE (

- Introduction.
- Mathematical model.
- Numerical model.
- Issues and Remedies.
 - Issue I: keep fractions bounded
 - Issue II: balance your "mass" account
 - Issue III: use a "compressible" PISO loop
- Conclusions.

Introduction



• Interested in gravity separation of two immiscible fluids.



- Use finite volume (OpenFOAM) code, adapted to our needs.
- Had and have some mass conservation issues.
- Found solutions to some but no all of them.

★ ∃ ▶ ★



- Drift Flux Model, Ishii & Hibiki (Two Fluid Model, Ishii 1987).
- The mixture is a single pseudo-fluid.
- Two phases: 1-continuous and 2-disperse.
- The phase fractions are bounded fields!
- Diferential Equations:
 - 1 mixture continuity
 - 2 mixture momentum
 - disperse phase transport
- Closure relation:
 - drift velocity model

ENIEF 2017 5 / 36

→ ∃ →



- Main fields:
 - α_2 : disperse phase volumetric fraction
 - **v**_m : center of mass (mixture) velocity
 - p_m : mixture pressure

- 4 @ > - 4 @ > - 4 @ >



- Main fields:
 - α_2 : disperse phase volumetric fraction
 - **v**_m : center of mass (mixture) velocity
 - p_m : mixture pressure

• Other important fields:

- \mathbf{v}_k : velocity of phase k = 1, 2
- $\alpha_1 = 1 \alpha_2$: continuous phase volumetric fraction
- $\rho_m = \alpha_1 \rho_1 + \alpha_2 \rho_2$: mixture density
- $\mathbf{j} = \alpha_1 \mathbf{v}_1 + \alpha_2 \mathbf{v}_2$: center of volume velocity
- $\mathbf{v}_{2j} = \mathbf{v}_2 \mathbf{j}$: "drift" velocity

ENIEF 2017 6 / 36

Differential Equations:

mixture continuity

$$\frac{\partial \rho_m}{\partial t} + \nabla \cdot (\rho_m \mathbf{v}_m) = 0$$



(日) (同) (三) (三)



Differential Equations:

Image: Image:

$$\frac{\partial \rho_m}{\partial t} + \nabla \cdot (\rho_m \mathbf{v}_m) = 0$$

2 mixture momentum

$$\frac{\partial \rho_m \mathbf{v}_m}{\partial t} + \nabla \cdot (\rho_m \mathbf{v}_m \mathbf{v}_m) = -\nabla \rho_m + \nabla \cdot (\tau + \tau^t) - \nabla \cdot \left[\frac{\alpha_2 \rho_1 \rho_2}{(1 - \alpha_2) \rho_m} \mathbf{v}_{2j} \mathbf{v}_{2j}\right] + \rho_m \mathbf{g}$$



▲ ■ ▶ ■ つへの ENIEF 2017 7 / 36

< ロ > < 同 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >

Differential Equations:

Image: Image:

$$\frac{\partial \rho_m}{\partial t} + \nabla \cdot (\rho_m \mathbf{v}_m) = 0$$

2 mixture momentum

$$\frac{\partial \rho_m \mathbf{v}_m}{\partial t} + \nabla \cdot (\rho_m \mathbf{v}_m \mathbf{v}_m) = -\nabla \rho_m + \nabla \cdot (\tau + \tau^t) - \nabla \cdot \left[\frac{\alpha_2 \rho_1 \rho_2}{(1 - \alpha_2) \rho_m} \mathbf{v}_{2j} \mathbf{v}_{2j}\right] + \rho_m \mathbf{g}$$

Isperse phase transport

$$\frac{\partial \alpha_2 \rho_2}{\partial t} + \nabla \cdot (\alpha_2 \rho_2 \mathbf{v}_m) = -\nabla \cdot \left(\frac{\alpha_2 \rho_1 \rho_2}{\rho_m} \mathbf{v}_{2j}\right)$$

Larreteguy - Caron - Barceló

ENIEF 2017 7 / 36

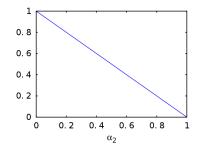
Closure relation:

drift velocity model (monodisperse)

$$\mathbf{v}_{2j} = \mathbf{V}_0(1 - lpha_2)$$

where

• $\mathbf{V}_0 = \frac{2}{9} \frac{g(\rho_1 - \rho_2)r_d^2}{\mu_1}$: terminal velocity of a droplet • $(1 - \alpha_2)$: factor to model the hindering effect



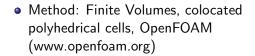


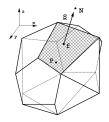
ENIEF 2017 8 / 36

★ ∃ >



Numerical model





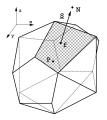
• • = • • = •



Numerical model



 Application: separationFoam (Larreteguy, Barceló, Caron), based on settlingFoam (Brennan, 2001)

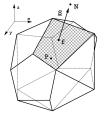




A = A = A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A

Numerical model

- Method: Finite Volumes, colocated polyhedrical cells, OpenFOAM (www.openfoam.org)
- Application: separationFoam (Larreteguy, Barceló, Caron), based on settlingFoam (Brennan, 2001)



- Features of separationFoam 3.2 (2017)
 - model for the mixture viscosity
 - thermal effects (i.e.: FWKO with fire tubes)
 - control system (i.e.: to control fire tubes power)
 - pressure gradient driven drift (i.e.: rotating flow separators)



Numerical model: Mixture continuity



• Finite volume approximation: \forall cell V^i with boundary ∂V^i

$$\int_{V^{i}} \frac{\partial \rho_{m}}{\partial t} dV + \int_{\partial V^{i}} (\mathbf{v}_{m} \rho_{m}) \cdot d\mathbf{S} = 0$$

dS: infinitesimal surface element.

• • = • •

Numerical model: Mixture continuity



• Finite volume approximation: \forall cell V^i with boundary ∂V^i

$$\int_{V^{i}} \frac{\partial \rho_{m}}{\partial t} dV + \int_{\partial V^{i}} (\mathbf{v}_{m} \rho_{m}) \cdot d\mathbf{S} = 0$$

dS : infinitesimal surface element.

• For cells bounded by flat faces

$$\frac{\partial \rho_m^i}{\partial t} V^i + \sum_f \Phi_m^f = 0$$

$$\begin{split} \Phi_m^f &= \rho_m^f \mathbf{v}_m^f \cdot \mathbf{S}^f : \text{ convective mass flux.} \\ (\cdot)^f : \text{ magnitude interpolated at the face.} \\ \mathbf{S}^f : \text{ face area vector.} \end{split}$$

ENIEF 2017 11 / 36

• • = • • = •



• Finite volume approximation:

$$\int_{V^{i}} \frac{\partial \alpha_{2}}{\partial t} dV + \int_{\partial V^{i}} (\mathbf{v}_{m} \alpha_{2}) \cdot d\mathbf{S} + \int_{\partial V^{i}} \left(\frac{\rho_{1}}{\rho_{m}} \mathbf{v}_{2j} \alpha_{2} \right) \cdot d\mathbf{S} = 0$$

→ ∃ →



• Finite volume approximation:

$$\int_{V^{i}} \frac{\partial \alpha_{2}}{\partial t} dV + \int_{\partial V^{i}} (\mathbf{v}_{m} \alpha_{2}) \cdot d\mathbf{S} + \int_{\partial V^{i}} \left(\frac{\rho_{1}}{\rho_{m}} \mathbf{v}_{2j} \alpha_{2} \right) \cdot d\mathbf{S} = 0$$

For flat faces

$$\frac{\partial \alpha_2^i}{\partial t} V^i + \sum_f \alpha_2^f \mathbf{v}_m^f \cdot \mathbf{S}^f + \sum_f \frac{\rho_1}{\rho_m^f} \alpha_2^f \mathbf{v}_{2j}^f \cdot \mathbf{S}^f = 0$$

Larreteguy - Caron - Barceló

ENIEF 2017 12 / 36

.



• Finite volume approximation:

$$\int_{V^{i}} \frac{\partial \alpha_{2}}{\partial t} dV + \int_{\partial V^{i}} (\mathbf{v}_{m} \alpha_{2}) \cdot d\mathbf{S} + \int_{\partial V^{i}} \left(\frac{\rho_{1}}{\rho_{m}} \mathbf{v}_{2j} \alpha_{2} \right) \cdot d\mathbf{S} = 0$$

• For flat faces

$$\frac{\partial \alpha_2^i}{\partial t} V^i + \sum_f \alpha_2^f \mathbf{v}_m^f \cdot \mathbf{S}^f + \sum_f \frac{\rho_1}{\rho_m^f} \alpha_2^f \mathbf{v}_{2j}^f \cdot \mathbf{S}^f = \mathbf{0}$$

• For consistency, we want to use the mass flux Φ_m^f in this equation.



• To do so, we define the disperse mass fraction $\tilde{\alpha}_2 = \frac{\rho_2}{\rho_m} \alpha_2$.

ENIEF 2017 13 / 36



- To do so, we define the disperse mass fraction $\tilde{\alpha}_2 = \frac{\rho_2}{\rho_m} \alpha_2$.
- We write the disperse phase mass transport equation as

$$\frac{\partial \rho_m^i \tilde{\alpha}_2^i}{\partial t} V^i + \sum_f \Phi_{\tilde{\alpha}}^f \tilde{\alpha}_2^f = 0$$

$$\begin{split} \Phi^{f}_{\tilde{\alpha}} &= \Phi^{f}_{m} + \Phi^{f}_{d} : \text{ total mass flux } \\ \Phi^{f}_{d} &= \rho_{1} \mathbf{v}^{f}_{2j} \cdot \mathbf{S}^{f} : \text{ drift mass flux. } \end{split}$$

ENIEF 2017 13 / 36

- **4 ∃ ≻** 4

Numerical model: Solution procedure



- Standard solution procedure for \mathbf{v}_m and p (PISO loop).
- $\tilde{\alpha}_2$ obtained from the disperse phase transport equation, which is not as simple as it might seem.

.



Problems with the original settlingFoam applied to water-oil separators:

• Symptom: most of the oil mass is lost!

¹A bounded downwind-upwind semi-discrete scheme for finite volume methods for phase separation problems, Axel E. Larreteguy, Luis F. Barceló, Pablo A. Caron, 1/Oct/2017, App. Math. Mod., Vol 50, pp 118-134; doi:10.1016/j.apm.2017.05.003

Larreteguy - Caron - Barceló

Time derivative in drift-flux method



Problems with the original settlingFoam applied to water-oil separators:

- Symptom: most of the oil mass is lost!
- Cause: the disperse phase keeps filling an already full cell, thus requiring truncation of disperse fraction to the range [0, 1].

¹A bounded downwind-upwind semi-discrete scheme for finite volume methods for phase separation problems, Axel E. Larreteguy, Luis F. Barceló, Pablo A. Caron, 1/Oct/2017, App. Math. Mod., Vol 50, pp 118-134; doi:10.1016/j.apm.2017.05.003

Larreteguy - Caron - Barceló

Time derivative in drift-flux method



Problems with the original settlingFoam applied to water-oil separators:

- Symptom: most of the oil mass is lost!
- Cause: the disperse phase keeps filling an already full cell, thus requiring truncation of disperse fraction to the range [0, 1].
- Remedy: new interpolation scheme (*downwind-upwind*¹) for the drift mass flux $\Phi_d^f = \rho_1 \mathbf{v}_{2j}^f \cdot \mathbf{S}^f$

¹A bounded downwind-upwind semi-discrete scheme for finite volume methods for phase separation problems, Axel E. Larreteguy, Luis F. Barceló, Pablo A. Caron, 1/Oct/2017, App. Math. Mod., Vol 50, pp 118-134; doi:10.1016/j.apm.2017.05.003

Larreteguy - Caron - Barceló

Time derivative in drift-flux method



Problems with the original settlingFoam applied to water-oil separators:

- Symptom: most of the oil mass is lost!
- Cause: the disperse phase keeps filling an already full cell, thus requiring truncation of disperse fraction to the range [0, 1].
- Remedy: new interpolation scheme (downwind-upwind¹) for the drift mass flux $\Phi_d^f = \rho_1 \mathbf{v}_{2j}^f \cdot \mathbf{S}^f$

Not all of our problems were gone, however.

¹A bounded downwind-upwind semi-discrete scheme for finite volume methods for phase separation problems, Axel E. Larreteguy, Luis F. Barceló, Pablo A. Caron, 1/Oct/2017, App. Math. Mod., Vol 50, pp 118-134; doi:10.1016/j.apm.2017.05.093

Larreteguy - Caron - Barceló

Time derivative in drift-flux method



Because ..

- Symptom: we still have **slight** losses of disperse phase.
- Cause?
- Remedy?

We are using conservative fluxes and a bounded scheme, so... how can it be?



The problem is tracked down to the following:

• Let us recall the disperse phase mass transport equation:

$$\frac{\partial \rho_m^i \tilde{\alpha}_2^i}{\partial t} V^i + \sum_f \Phi_{\tilde{\alpha}}^f \tilde{\alpha}_2^f = 0$$

- The hidden problem: the time derivative is applied to the product of two variables that are not independent, for
 ρ_m = ρ₁ρ₂/(ρ₂ − (ρ₂ − ρ₁)α̃₂).
- In other words: the time derivative is applied to a non-linear function of $\tilde{\alpha}_2$
- We need to linearize it.

ENIEF 2017 18 / 36

• • = • • = •



19 / 36

Original settlingFoam algorithm for solving for $\tilde{\alpha}$ and ρ_m :

• The linearization is done by splitting the derivative of the product

$$\rho_{m}^{i} \frac{\partial \tilde{\alpha}_{2}^{i}}{\partial t} V^{i} + \tilde{\alpha}_{2}^{i} \frac{\partial \rho_{m}^{i}}{\partial t} V^{i} + \sum_{f} \Phi_{\tilde{\alpha}}^{f} \tilde{\alpha}_{2}^{f} = 0$$

and neglecting the second term (i.e.: consider $\rho_m^i = \rho_m^{io} = const$).

イロト イポト イヨト イヨト



Original settlingFoam algorithm for solving for $\tilde{\alpha}$ and ρ_m :

• The linearization is done by splitting the derivative of the product

$$\rho_{m}^{i} \frac{\partial \tilde{\alpha}_{2}^{i}}{\partial t} V^{i} + \tilde{\alpha}_{2}^{i} \frac{\partial \rho_{m}^{i}}{\partial t} V^{i} + \sum_{f} \Phi_{\tilde{\alpha}}^{f} \tilde{\alpha}_{2}^{f} = 0$$

and neglecting the second term (i.e.: consider $\rho_m^i = \rho_m^{io} = const$).

- The solution algorithm is the following:
 - **1** solve equation to get the updated mass fraction $\tilde{\alpha}_2$;
 - 2 evaluate the updated density as $\rho_m = \rho_1 + (\rho_1/\rho_2 1)\tilde{\alpha}_2$; and
 - **③** evaluate the updated disperse volume fraction $\alpha_2 = \rho_m \tilde{\alpha}_2 / \rho_2$.

(日) (同) (日) (日) (日)



We propose the following

Predictor-Corrector algorithm:

• obtain a *predicted* $\tilde{\alpha}_2^{i*}$ by aplying step i) of the original algorithm,

• • = • •



We propose the following

Predictor-Corrector algorithm:

- **(**) obtain a *predicted* $\tilde{\alpha}_2^{i*}$ by aplying step i) of the original algorithm,
- 2 compute the updated mass as $m_2^i = \rho_m^{io} \tilde{\alpha}_2^{i*} V^i$,

< 回 ト < 三 ト < 三 ト



We propose the following

Predictor-Corrector algorithm:

- In obtain a predicted $\tilde{\alpha}_2^{i*}$ by aplying step i) of the original algorithm,
- ② compute the updated mass as $m_2^i=
 ho_m^{io} ildelpha_2^{i*}V^i$,
- **③** compute the updated volume fraction as $\alpha_2^i = m_2^i / (\rho_2 V^i)$,

▲圖 ▶ ▲ 国 ▶ ▲ 国 ▶



We propose the following

Predictor-Corrector algorithm:

- () obtain a *predicted* $\tilde{\alpha}_2^{i*}$ by aplying step i) of the original algorithm,
- ② compute the updated mass as $m_2^i =
 ho_m^{io} ilde lpha_2^{i*} V^i$,
- ${f 3}$ compute the updated volume fraction as $lpha_2^i=m_2^i/(
 ho_2V^i)$,
- update the mixture density as $ho_m^i = (1 lpha_2^i)
 ho_1 + lpha_2^i
 ho_2$, and finally

イロト イヨト イヨト



We propose the following

Predictor-Corrector algorithm:

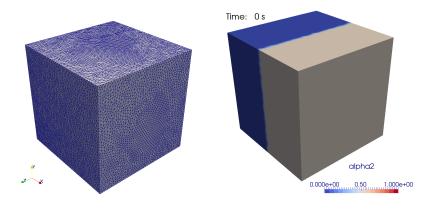
- In obtain a predicted $\tilde{\alpha}_2^{i*}$ by aplying step i) of the original algorithm,
- ② compute the updated mass as $m_2^i = \rho_m^{io} \tilde{lpha}_2^{i*} V^i$,
- ③ compute the updated volume fraction as $lpha_2^i=m_2^i/(
 ho_2V^i)$,
- ④ update the mixture density as $ho_m^i = (1-lpha_2^i)
 ho_1 + lpha_2^i
 ho_2$, and finally
- **5** update the mass fraction as $\tilde{\alpha}_2 = \rho_2 \alpha_2 / \rho_m$.

ENIEF 2017 20 / 36

イロト イヨト イヨト



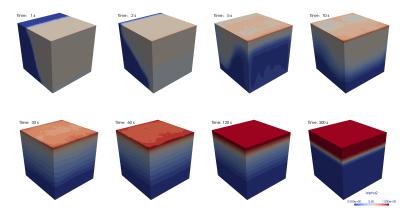
Predictor-corrector algorithm: test - density driven stirring and separation of a water-oil mixture in a cube (in a *horrible* mesh)



- **→ →** •



Hard test in a *horrible* mesh: results



Unit cube - snapshots of volume fraction at selected times

Larreteguy - Caron - Barceló

Time derivative in drift-flux method

ENIEF 2017 22 / 36

(日) (同) (三) (三)



• The proposed predictor-corrector algorithm has been found to conserve mass up to machine precision *as long as there is no need for truncation of the disperse phase fraction field.*



- The proposed predictor-corrector algorithm has been found to conserve mass up to machine precision *as long as there is no need for truncation of the disperse phase fraction field.*
- But, what about accuracy?

A = A = A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A



The bottle experiment:

▶ < ≧ ▶ ≧ ∽ < < ENIEF 2017 24 / 36

(日) (同) (三) (三)



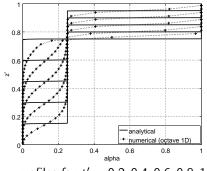
The *bottle experiment*: an academic 1D verification example

- Uniform mixture in a tank, left to separate (height H = 10m).
- $\mathbf{j} = \mathbf{0} \Rightarrow$ mixture mass and momentum equations not requiered.
- Analytical solution available.
- 1D numerical solution also available (sim1D: 1D Octave version of separationFoam).
- Case: oil in water ($ho_1=1000 kg/m^3,
 ho_1=900 kg/m^3$)
- Discretization details for both separationFoam and sim1D
 - uniform spatial discretization (40 cells)
 - fixed time step ($\Delta t = \Delta z/V_0 \Rightarrow \mathit{Co}_{drift} = 1$).

イロト 不得下 イヨト イヨト 二日



The bottle experiment: analytical and sim1D solutions.



 α_2 profiles for t' = 0.2, 0.4, 0.6, 0.8, 1.0

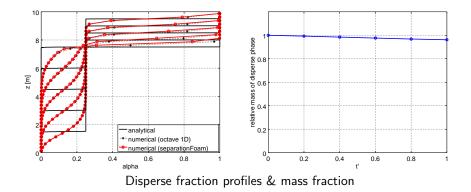
Dimensionless solution (valid for all mixtures, drop size and tank height).

Larreteguy - Caron - Barceló

ENIEF 2017 26 / 36



The bottle experiment (oil in water): original algorithm



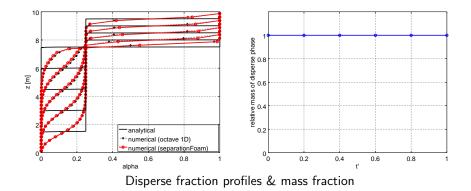
Mass NOT conserved (despite that fraction truncation is not required).

Larreteguy - Caron - Barceló

ENIEF 2017 27 / 36



The bottle experiment (oil in water): predictor-corrector algorithm



Mass IS conserved (as long as fractions are not truncated).

Larreteguy - Caron - Barceló

Time derivative in drift-flux method

ENIEF 2017 2

28 / 36



We seem to have solved Issue II: "we still have slight losses of disperse phase"

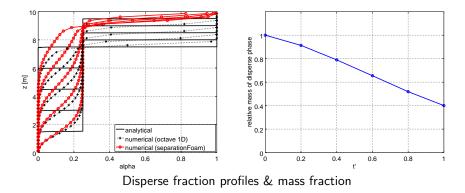
- Cause: bad accounting of mass inventory.
- Remedy: predictor-corrector algorithm.

But, is it really so?

Let us consider "air" in water ($\rho_1 = 1000 kg/m^3$, $\rho_2 = 1 kg/m^3$).



The bottle experiment (air in water): original algorithm



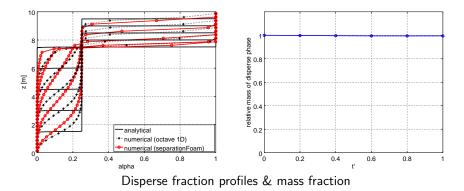
- Huge mass loss (due to both bad mass accounting AND truncation).
- separationFoam and sim1D solutions differ (they shouldn't!)

Larreteguy - Caron - Barceló

ENIEF 2017 30 / 36



The bottle experiment (air in water): predictor-corrector algorithm



- Mass IS conserved, but...
- separationFoam and sim1D solutions still differ!

Larreteguy - Caron - Barceló

Time derivative in drift-flux method

ENIEF 2017 31 / 36



So, we are still in trouble \Rightarrow **Issue III**.

Hint: It has to be related to the equations **not solved in** sim1D (i.e. those of the PISO loop).

• Symptom: separationFoam and sim1D solutions differ.

→ Ξ →



So, we are still in trouble \Rightarrow **Issue III**.

Hint: It has to be related to the equations **not solved in** sim1D (i.e. those of the PISO loop).

- Symptom: separationFoam and sim1D solutions differ.
- Cause: PISO loop is not dealing well with the "compressible" mixture

$$\sum_{f} \left(\frac{1}{a_{i}}\right)_{f} (\nabla p)_{f} \cdot \mathbf{S}^{f} = \underbrace{\frac{\partial \rho_{m}^{i}}{\partial t} V^{i}}_{\text{hum the ded with which}} + \sum_{f} \Phi_{U}^{f}$$

how to deal with this?

leading to a solution that has $\mathbf{j} \neq \mathbf{0}$.

くほと くほと くほと



So, we are still in trouble \Rightarrow **Issue III**.

Hint: It has to be related to the equations **not solved in** sim1D (i.e. those of the PISO loop).

- Symptom: separationFoam and sim1D solutions differ.
- Cause: PISO loop is not dealing well with the "compressible" mixture

$$\sum_{f} \left(\frac{1}{a_{i}}\right)_{f} (\nabla p)_{f} \cdot \mathbf{S}^{f} = \underbrace{\frac{\partial \rho_{m}^{i}}{\partial t} V^{i}}_{\text{how to deal with this?}} + \sum_{f} \Phi_{U}^{f}$$

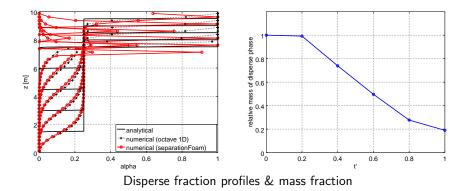
leading to a solution that has $\mathbf{j} \neq \mathbf{0}$.

• Remedy: a possible approach is to use $\partial \rho_m^i / \partial t$ from the previous step.

(日) (同) (日) (日) (日)



The bottle experiment (air in water): predictor-corrector algorithm



The water front from separationFoam and sim1D are much closer,but the air front is a total disaster!

Larreteguy - Caron - Barceló

Time derivative in drift-flux method

ENIEF 2017 33 / 36



To conclude this issue:

• The "compressible" PISO explains the differences between the two numerical solutions, and is the right way to go, although some problems remain.



To conclude this issue:

- The "compressible" PISO explains the differences between the two numerical solutions, and is the right way to go, although some problems remain.
- We have identified that the problems are related to the way OpenFOAM imposes a reference pressure at some point, which is required for closed domains to overcome the "up to an aditive constant" indefinition of the pressure field.

• • = • • = •



To conclude this issue:

- The "compressible" PISO explains the differences between the two numerical solutions, and is the right way to go, although some problems remain.
- We have identified that the problems are related to the way OpenFOAM imposes a reference pressure at some point, which is required for closed domains to overcome the "up to an aditive constant" indefinition of the pressure field.
- We tracked down the problem to either inadequate use of the available boundary conditions (meaning US) or inadequate handling of these boundary conditions in the coding (meaning THEM).

(日) (周) (三) (三)



To conclude this issue:

- The "compressible" PISO explains the differences between the two numerical solutions, and is the right way to go, although some problems remain.
- We have identified that the problems are related to the way OpenFOAM imposes a reference pressure at some point, which is required for closed domains to overcome the "up to an aditive constant" indefinition of the pressure field.
- We tracked down the problem to either inadequate use of the available boundary conditions (meaning US) or inadequate handling of these boundary conditions in the coding (meaning THEM).
- We owe you the final answer... Still working on it.

(日) (周) (三) (三)

Conclusions



- We have found some issues and developed and implemented remedies for a drift flux model as applied to separation problems.
- The resulting code is able to give solutions even with extreme density ratios (i.e. 1/1000) and high Courant numbers.
- There are unsolved issues related to the use and/or coding of boundary conditions and their interaction with the "compressible" PISO loop.

A D A D A D A

Acknowledgments



The authors wish to thank

- Universidad Argentina de la Empresa (support under Grant P17T03)
- ENIEF 2017 organizers
- You, for your kind attention

★ ∃ ▶ ★